

## NUMERICAL STUDY FOR A SECOND ORDER FREDHOLM INTEGRO-DIFFERENTIAL EQUATION BY APPLYING GALERKIN-Chebyshev-WAVELETS METHOD

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**Abstract.** In the present paper, we apply the Galerkin method using Chebyshev wavelets to approximate the exact solution for a second order Fredholm integro-differential equation with initial conditions. This numerical method gives us a nonlinear algebraic system that would be solved using the Picard successive approximations technique. Furthermore, we show the validity and the ability of the proposed method through some illustrative examples.

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### 1. Introduction

Recently, there has been a growing interest in the Integro-Differential Equations (IDEs), for these kinds of equations can be found in modeling real phenomena in many fields of sciences, physics, chemistry, biology and engineering problems, such as epidemic models [1, 2], the Boltzmann kinetic equation [3], and the Vlasov and Landau equations [4]. However, the exact solution to such equations is usually difficult to obtain, so the researchers use different numerical methods to approach the exact solution such as the Pade approximation, the Legendre-Galerkin method [5], the Hermite wavelet [6], the Haar wavelet [7], the Chebyshev wavelet collocation method [8, 9], the wavelet-Galerkin method [10], the Laguerre wavelets collocation method [11], the Laplace decomposition method [12], Bernoulli polynomials [13], and the B-spline method [14, 15].

In this article, we consider the following Fredholm integro-differential equation with initial conditions:

$$\begin{cases} \phi(x) = f(x) + \int_0^1 K(x, y, \phi(y), \phi'(y), \phi''(y)) dy, \\ \phi(0) = \rho_1, \quad \phi'(0) = \rho_2, \end{cases} \quad (1)$$

where  $\phi(x), f(x) \in H^2([0, 1])$ ,  $K, \partial_x K, \partial_x^2 K \in C([0, 1]^2 \times \mathbb{R}^3)$ , and  $\rho_1, \rho_2 \in \mathbb{R}$ .

We mention that this type of equation has a specific form, whereas the unknown function and its derivatives appear inside of the nonlinear kernel, so the main result of this paper is to introduce a numerical method to approach the exact solution of the equation (1) by applying the Galerkin method with Chebyshev-Wavelets.

This paper is structured as follows: In section 2, we introduce some Chebyshev wavelet properties and function approximation. In section 3, we build the integration operational matrices. In section 4, we explain our numerical method, and in section 5, we prove the convergence analysis of the proposed method. Finally, in section 6, we demonstrate the accuracy and efficiency of our method with some illustrative examples.

## 2. Properties of Chebyshev wavelets

### 2.1. Chebyshev polynomials

The Chebyshev polynomials are obtained by expanding the formula:

$$T_n(x) = \cos(n \arccos(x)),$$

and satisfy the following recurrence relation:

$$\begin{cases} T_{n+1}(x) = 2xT_n(x) - T_{n-1}(x), \\ T_1(x) = x, \quad T_0(x) = 1, \quad \text{for } n \geq 1. \end{cases}$$

These polynomials are orthogonal with respect to the weight function  $w(x) = \frac{1}{\sqrt{1-x^2}}$  on  $[-1, 1]$ , and

$$\int_{-1}^1 w(x) T_n(x) T_m(x) dx = \begin{cases} \pi, & n = m = 0, \\ \pi/2, & n = m \neq 0, \\ 0, & n \neq m. \end{cases}$$

The first few Chebyshev polynomials are:

$$\begin{aligned} T_0(x) &= 1, \\ T_1(x) &= x, \\ T_2(x) &= 2x^2 - 1, \\ T_3(x) &= 4x^3 - 3x, \\ T_4(x) &= 8x^4 - 8x^2 + 1, \\ T_5(x) &= 16x^5 - 20x^3 + 5x. \end{aligned}$$

## 2.2. Chebyshev wavelets

The wavelets are stretched versions of the original wavelet  $T$  with the same basic shape but at a different scale and frequency. We construct them from the translation parameter  $b$  and the scaling parameter  $a$ , then we have the following family of wavelets:

$$\theta_{a,b}(x) = \frac{1}{\sqrt{|a|}} T\left(\frac{x-b}{a}\right), \quad a, b \in \mathbb{R}, \quad a \neq 0.$$

So, the Chebyshev wavelets are defined as:

$$\theta_{i,j}(x) = \begin{cases} 2^{\frac{k}{2}} \tilde{T}_j(2^k x - 2i + 1), & \frac{i-1}{2^{k-1}} \leq x \leq \frac{i}{2^{k-1}}, \\ 0, & \text{otherwise,} \end{cases}$$

and

$$\tilde{T}_j(t) = \begin{cases} \frac{1}{\sqrt{\pi}}, & j = 0, \\ \sqrt{\frac{2}{\pi}} T_j(t), & j > 0, \end{cases}$$

where  $k$  is a positive integer number,  $j = 0, 1, 2, \dots, n-1$ ,  $i = 1, 2, \dots, 2^{k-1}$  and  $T_j$  is the Chebyshev polynomial of degree  $j$ . However, the family of Chebyshev wavelets  $\{\theta_{i,j}\}$  defines an orthonormal basis for  $L_{w_k}^2([0, 1])$  such that:

$$w_k(x) = \begin{cases} w_{1,k}(x), & 0 \leq x < \frac{1}{2^{k-1}}, \\ w_{2,k}(x), & \frac{1}{2^{k-1}} \leq x < \frac{2}{2^{k-1}}, \\ \vdots & \vdots \\ w_{2^{k-1},k}(x), & \frac{2^{k-1}-1}{2^{k-1}} \leq x < 1, \end{cases}$$

where  $w_{i,k}(x) = w(2^k x - 2i + 1)$ . The Chebyshev wavelets charts are proven for  $k = 1$  and  $n = 1, 2, 3, 4$  in Figure 1.

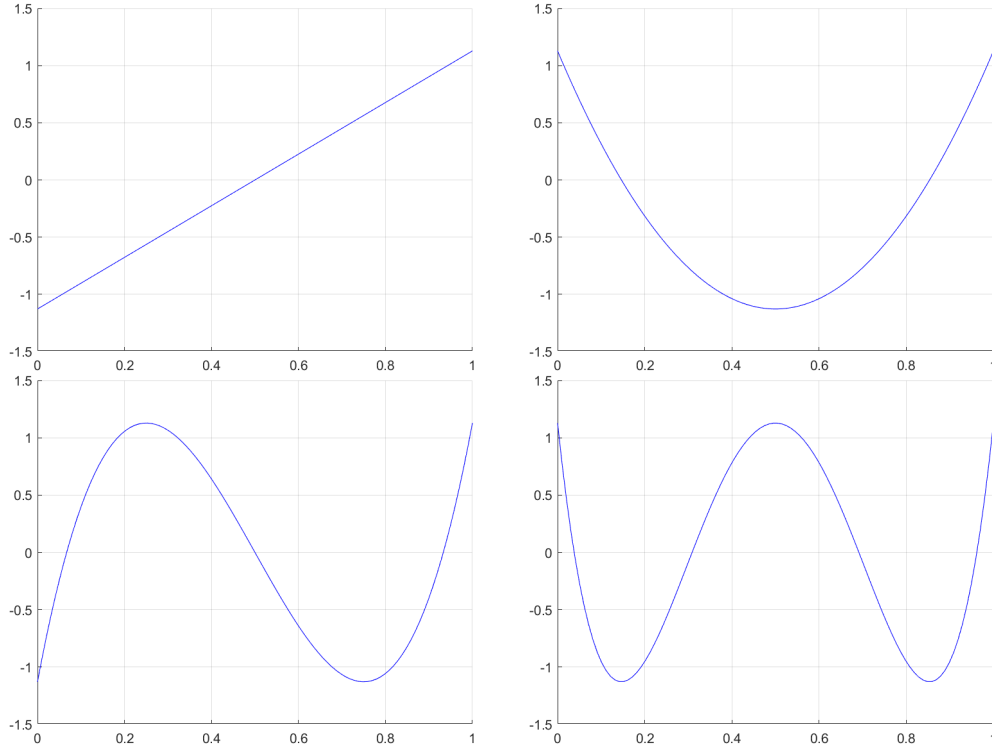


Fig. 1. Chebyshev wavelets for  $k = 1$

For any function  $\phi(x)$  in  $L^2_{w_k}([0, 1])$  can be expanded as:

$$\phi(x) = \sum_{i=1}^{\infty} \sum_{j=0}^{\infty} c_{i,j} \theta_{i,j}(x), \quad (2)$$

where  $c_{i,j} = \langle \phi, \theta_{i,j} \rangle$ , such that  $\langle \cdot, \cdot \rangle$  is the inner product in  $L^2_{w_k}([0, 1])$ . So we approximate the function  $\phi(x)$  by truncating the infinite series (2):

$$\phi_n(x) = \sum_{i=1}^{2^{k-1}} \sum_{j=0}^{n-1} c_{i,j} \theta_{i,j}(x) = C^T P(x), \quad (3)$$

where,  $C^T$  and  $P(x)$  are  $2^{k-1}n \times 1$  matrices:

$$C^T = [c_{1,0}, c_{1,1}, \dots, c_{1,n-1}, c_{2,0}, c_{2,1}, \dots, c_{2,n-1}, \dots, c_{2^{k-1},0}, \dots, c_{2^{k-1},n-1}],$$

and

$$P(x) = [\theta_{1,0}, \theta_{1,1}, \dots, \theta_{1,n-1}, \theta_{2,0}, \theta_{2,1}, \dots, \theta_{2,n-1}, \dots, \theta_{2^{k-1},0}, \theta_{2^{k-1},1}, \dots, \theta_{2^{k-1},n-1}]^T.$$

### 3. Chebyshev wavelet operational matrix of integration

If we take  $k = 2$ , then the matrices  $C^T$  and  $P(x)$  would be:

$$C^T = [c_{1,0}, c_{1,1}, \dots, c_{1,n-1}, c_{2,0}, c_{2,1}, \dots, c_{2,n-1}].$$

$$P(x) = [\theta_{1,0}(x), \theta_{1,1}(x), \dots, \theta_{1,n-1}(x), \theta_{2,0}(x), \theta_{2,1}(x), \dots, \theta_{2,n-1}(x)].$$

Let  $W_n$  be a matrix that contains the coefficients of Chebyshev wavelets:

$$W_n = \begin{pmatrix} F_n & O_n \\ O_n & \tilde{F}_n \end{pmatrix},$$

$$F_n = \frac{2}{\sqrt{\pi}} \begin{pmatrix} 1 & -\sqrt{2} & \dots & (-1)^{n-1}\sqrt{2} \\ 0 & 4\sqrt{2} & \dots & \vdots \\ \vdots & \vdots & \ddots & \\ 0 & 0 & \dots & 4^{n-1}\sqrt{2} \end{pmatrix}, \quad \tilde{F}_n = \frac{2}{\sqrt{\pi}} \begin{pmatrix} 1 & -3\sqrt{2} & \dots & T_{n-1}(-3)\sqrt{2} \\ 0 & 4\sqrt{2} & \dots & \vdots \\ \vdots & \vdots & \ddots & \\ 0 & 0 & \dots & 4^{n-1}\sqrt{2} \end{pmatrix},$$

and

$$X_n(x) = (1, x, x^2, \dots, x^{n-1}, 1, x, x^2, \dots, x^{n-1}),$$

$$P_n(x) = [\theta_{1,0}, \theta_{1,1}, \dots, \theta_{1,n-1}, \theta_{2,0}, \theta_{2,1}, \dots, \theta_{2,n-1}].$$

So, we can write:

$$P_n(x) = X_n(x)W_n.$$

Let  $N_n$  be an integral matrix in classical basis for polynomial space:

$$N_n = \begin{pmatrix} G_n & O_n \\ O_n & \tilde{G}_n \end{pmatrix},$$

$$G_n = \begin{pmatrix} 0 & 1 & 0 & 0 & \dots & 0 \\ 0 & 0 & \frac{1}{2} & 0 & \dots & 0 \\ 0 & 0 & 0 & \frac{1}{3} & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & \frac{1}{n} \end{pmatrix}, \quad \tilde{G}_n = \begin{pmatrix} -\frac{1}{2} & 1 & 0 & 0 & \dots & 0 \\ -\frac{1}{2^2} & 0 & \frac{1}{2} & 0 & \dots & 0 \\ 0 & 0 & 0 & \frac{1}{3} & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ -\frac{1}{2^n} & 0 & 0 & 0 & \dots & \frac{1}{n} \end{pmatrix},$$

then, we have the first integration matrix:

$$\int_0^x C^T P_n(y) dy = C^T M_n P_{n+1}(x) = C^T Q_1(x),$$

and the double integration matrix:

$$\int_0^x \int_0^z C^T P_n(y) dy dz = C^T M_n M_{n+1} P_{n+2}(x) = C^T Q_2(x),$$

where,

$$M_n = W_n^{-1} N_n W_{n+1}.$$

#### 4. Description of method

Let the following Fredholm integro-differential equation:

$$\begin{cases} \phi(x) = f(x) + \int_0^1 K(x, y, \phi(y), \phi'(y), \phi''(y)) dy, \\ \phi(0) = \rho_1, \quad \phi'(0) = \rho_2. \end{cases} \quad (4)$$

First, we derive the equation (4) twice to obtain the following equation:

$$\phi''(x) = f''(x) + \int_0^1 \partial_x^2 K(x, y, \phi(y), \phi'(y), \phi''(y)) dy, \quad (5)$$

we approach the function  $\phi''(x)$  using Chebyshev wavelets as:

$$\phi_n''(x) = C^T P_n(x), \quad (6)$$

integrate the equation (6) from 0 to  $x$ , we obtain:

$$\phi_n'(x) = \rho_2 + C^T Q_1(x), \quad (7)$$

integrate again (7) from 0 to  $x$ , then we get:

$$\phi_n(x) = \rho_1 + \rho_2 x + C^T Q_2(x), \quad (8)$$

by substituting (6), (7) and (8) in (5), we obtain:

$$C^T P_n(x) = f''(x) + \int_0^1 \partial_x^2 K(x, y, \rho_1 + \rho_2 y + C^T Q_2(y), \rho_2 + C^T Q_1(y), C^T P_n(y)) dy. \quad (9)$$

Multiplying equation (9) by  $\theta_{i,j}(x) w_{i,2}(x)$  for  $i = 1, 2$  and  $j = 0, 1, \dots, n - 1$ , after that we integrate with respect to  $x$  from 0 to 1, then we get the following nonlinear

algebraic system:

$$c_{i,j} = y_{i,j} + \int_0^1 \int_0^1 \theta_{i,j}(x) w_{i,2}(x) \partial_x^2 K(x, y, \rho_1 + \rho_2 y + C^T Q_2(y), \rho_2 + C^T Q_1(y), C^T P_n(y)) dy dx, \quad (10)$$

where  $y_{i,j} = \langle f'', \theta_{i,j} \rangle$ . However, we can find the vector  $C^T$  solution of the above system (10) by using the Picard successive approximations method, then we substitute them in (8) to get the numerical solution of the proposed equation (1).

## 5. The convergence analysis

First, to prove the convergence analysis of the proposed numerical process described above, consider the Sobolev space  $\mathcal{H} = H^2([0, 1], \mathbb{R})$  equipped with the following norm:

$$\forall \phi \in \mathcal{H}, \|\phi\|_{\mathcal{H}} = \|\phi\|_{L^2[0,1]} + \|\phi'\|_{L^2[0,1]} + \|\phi''\|_{L^2[0,1]}$$

Furthermore, let's suppose the following assumptions:

$$(\mathcal{A}) \left\{ \begin{array}{l} \exists A_m, B_m, C_m > 0, \text{ where } m = 0, 1, 2. \forall x, y \in [0, 1], \forall u, \bar{u}, v, \bar{v}, w, \bar{w} \in \mathbb{R}, \\ |K(x, y, u, v, w) - K(x, y, \bar{u}, \bar{v}, \bar{w})| \leq A_0 |u - \bar{u}| + B_0 |v - \bar{v}| + C_0 |w - \bar{w}|, \\ |\partial_x K(x, y, u, v, w) - \partial_x K(x, y, \bar{u}, \bar{v}, \bar{w})| \leq A_1 |u - \bar{u}| + B_1 |v - \bar{v}| + C_1 |w - \bar{w}|, \\ |\partial_x^2 K(x, y, u, v, w) - \partial_x^2 K(x, y, \bar{u}, \bar{v}, \bar{w})| \leq A_2 |u - \bar{u}| + B_2 |v - \bar{v}| + C_2 |w - \bar{w}|, \\ 0 < \gamma = \max \left\{ \sum_{m=0}^2 A_m, \sum_{m=0}^2 B_m, \sum_{m=0}^2 C_m \right\} < 1. \end{array} \right.$$

**Theorem 1** According to assumptions  $(\mathcal{A})$ , the approximate solution  $\phi_n$  converges to the exact solution  $\phi$  in the space  $\mathcal{H}$ .  $\square$

PROOF Consider the following operator  $T$  defined from  $\mathcal{H}$  to itself by:

$$\forall x \in [0, 1], \quad T : \mathcal{H} \longrightarrow \mathcal{H}$$

$$\phi \longmapsto T(\phi)(x) = f(x) + \int_0^1 K(x, y, \phi(y), \phi'(y), \phi''(y)) dy.$$

So, the exact solution  $\phi$  of our equation (1) and its derivatives  $\phi'$  and  $\phi''$  can be represented by the following system:

$$\begin{cases} \phi(x) = T(\phi)(x), \\ \phi'(x) = T'(\phi)(x), \\ \phi''(x) = T''(\phi)(x). \end{cases}$$

By applying the Galerkin projection method using Chebyshev wavelets given by (3),

the previous system will be approached by:

$$\begin{cases} \phi_n(x) = T_n(\phi_n)(x), \\ \phi_n'(x) = T_n'(\phi_n)(x), \\ \phi_n''(x) = T_n''(\phi_n)(x). \end{cases}$$

It's clear that by applying the triangle inequality, we get:

$$\begin{aligned} |\phi_n(x) - \phi(x)| &= |T_n(\phi_n) - T(\phi)| = |T_n(\phi_n) - T(\phi_n) + T(\phi_n) - T(\phi)|, \\ &\leq |T_n(\phi_n) - T(\phi_n)| + |T(\phi_n) - T(\phi)|. \end{aligned}$$

According to assumptions ( $\mathcal{A}$ ) and Cauchy Schwarz inequality, we achieve:

$$\begin{aligned} |T(\phi_n) - T(\phi)| &= \left| \int_0^1 (K(x, y, \phi_n(y), \phi_n'(y), \phi_n''(y)) - K(x, y, \phi(y), \phi'(y), \phi''(y))) dy \right|, \\ &\leq A_0 \int_0^1 |\phi_n(y) - \phi(y)| dy + B_0 \int_0^1 |\phi_n'(y) - \phi'(y)| dy + C_0 \int_0^1 |\phi_n''(y) - \phi''(y)| dy, \\ &\leq A_0 \|\phi_n - \phi\|_{L^2[0,1]} + B_0 \|\phi_n' - \phi'\|_{L^2[0,1]} + C_0 \|\phi_n'' - \phi''\|_{L^2[0,1]}. \end{aligned} \quad (11)$$

On the other hand, the study [16] assumes that sequence  $S_n = T_n(\phi_n)$  is convergent, and gives us the following error convergence order:

$$|T_n(\phi_n) - T(\phi_n)| \leq \mathcal{O}(n^{\mu_0}) \rightarrow 0. \quad (12)$$

So, from inequalities (11) and (12), we obtain:

$$\|\phi_n - \phi\|_{L^2[0,1]} \leq A_0 \|\phi_n - \phi\|_{L^2[0,1]} + B_0 \|\phi_n' - \phi'\|_{L^2[0,1]} + C_0 \|\phi_n'' - \phi''\|_{L^2[0,1]} + \mathcal{O}(n^{\mu_0}).$$

Similarly, we find:

$$\begin{aligned} \|\phi_n' - \phi'\|_{L^2[0,1]} &\leq A_1 \|\phi_n - \phi\|_{L^2[0,1]} + B_1 \|\phi_n' - \phi'\|_{L^2[0,1]} + C_1 \|\phi_n'' - \phi''\|_{L^2[0,1]} + \mathcal{O}(n^{\mu_1}), \\ \|\phi_n'' - \phi''\|_{L^2[0,1]} &\leq A_2 \|\phi_n - \phi\|_{L^2[0,1]} + B_2 \|\phi_n' - \phi'\|_{L^2[0,1]} + C_2 \|\phi_n'' - \phi''\|_{L^2[0,1]} + \mathcal{O}(n^{\mu_2}). \end{aligned}$$

Therefore, if we put  $\mathcal{O}(n^\mu) = \mathcal{O}(n^{\mu_0}) + \mathcal{O}(n^{\mu_1}) + \mathcal{O}(n^{\mu_2}) \rightarrow 0$ , and according to  $0 < \gamma < 1$ , we can conclude that:

$$\|\phi_n - \phi\|_{\mathcal{H}} \leq \frac{\mathcal{O}(n^\mu)}{1 - \gamma} \rightarrow 0.$$

Which confirms the convergence of the approximate solution  $\phi_n$  to  $\phi$  in the space  $\mathcal{H}$ . ■

## 6. Examples

In this section, we apply our numerical method on two illustrative examples, in order to verify its accuracy and validity. For this reason, we define the error function



$E_n$  by:

$$E_n = \|\phi_n - \phi\|_{\mathcal{H}},$$

where  $\phi$  is the exact solution of the example, and  $\phi_n$  is the numerical solution obtained using Matlab computation software.

### 6.1. Example 1

Consider the following equation:

$$\begin{cases} \phi(x) = f(x) + \int_0^1 K(x, y, \phi(y), \phi'(y), \phi''(y)) dy, & \forall x \in [0, 1], \\ \phi(0) = 0, \quad \phi'(0) = 0, \end{cases}$$

with

$$f(x) = x \sin(x) + \frac{1}{2} \sin(x)(2\ln(2) - 1),$$

$$K(x, y, u, v, w) = -\frac{1}{2} \sin(x) \ln[1 + \sin(y)u + \cos(y)v - \frac{1}{2} \sin(y)(w - \sin(y))],$$

and the exact solution is  $\phi(x) = x \sin(x)$ .

### 6.2. Example 2

Let the Fredholm integro-differential equation:

$$\begin{cases} \phi(x) = f(x) - \int_0^1 \frac{3}{4} \cos(x) \sin \left[ \frac{1}{2} e^{-y} (2\phi''(y) - \phi'(y) - \phi(y)) \right] ds, & \forall x \in [0, 1], \\ \phi(0) = \frac{1}{4}, \quad \phi'(0) = -\frac{3}{4}, \end{cases}$$

with

$$f(x) = \begin{cases} \frac{1}{4} \left[ (2x-1)^2 e^x + \cos(x) \left[ \cos\left(\frac{1}{2}\right) - 2\cos(2) + \cos\left(\frac{7}{2}\right) \right] \right], & 0 \leq x \leq \frac{1}{2}, \\ \frac{1}{4} \left[ -(2x-1)^2 e^x + \cos(x) \left[ \cos\left(\frac{1}{2}\right) - 2\cos(2) + \cos\left(\frac{7}{2}\right) \right] \right], & \frac{1}{2} \leq x \leq 1, \end{cases}$$

and the exact solution is:

$$\phi(x) = \begin{cases} \frac{1}{4} (2x-1)^2 e^x, & 0 \leq x \leq \frac{1}{2}, \\ -\frac{1}{4} (2x-1)^2 e^x, & \frac{1}{2} \leq x \leq 1. \end{cases}$$

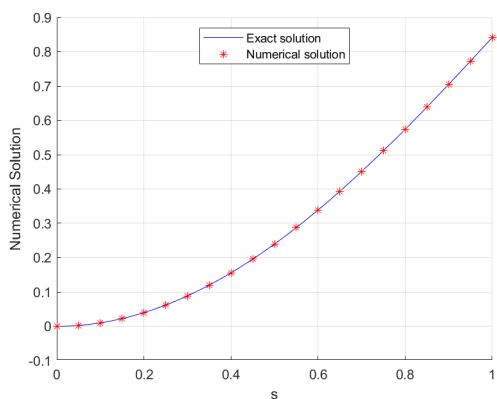


Fig. 2. Graphical comparison of  $\phi$  vs  $\phi_n$  for Example 1, with  $n = 7$

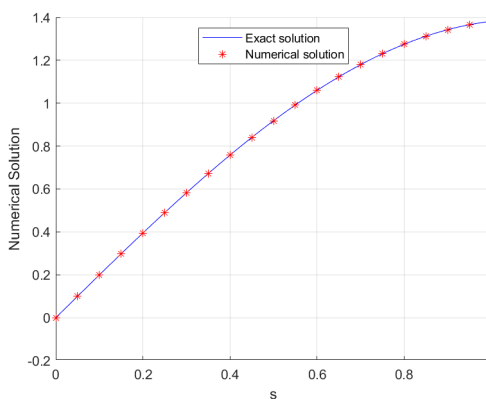


Fig. 3. Graphical comparison of  $\phi'$  vs  $\phi'_n$  for Example 1, with  $n = 7$

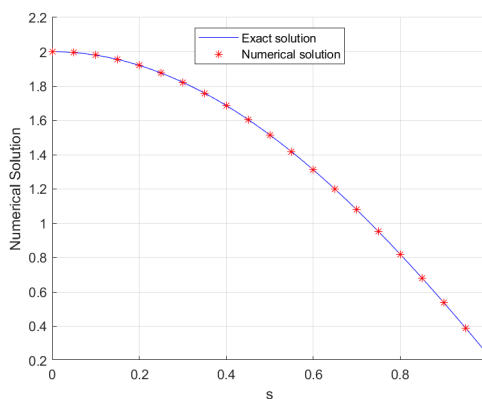


Fig. 4. Graphical comparison of  $\phi''$  vs  $\phi''_n$  for Example 1, with  $n = 7$

Table 1. Numerical results of Example 1

$n$	$E_n$	CPU time
3	6.9888e-05	0.109
4	1.4336e-06	0.123
5	4.9340e-08	0.154
6	6.7237e-10	0.194
7	5.4441e-10	0.310

Table 2. Numerical results of Example 2

$n$	$E_n$	CPU time
3	7.9499e-04	0.081
4	2.7590e-05	0.108
5	7.6178e-07	0.129
6	1.7419e-08	0.169
7	8.5031e-09	0.372

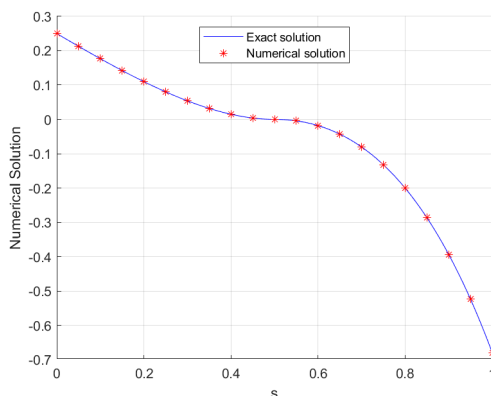


Fig. 5. Graphical comparison of  $\phi$  vs  $\phi_n$  for Example 2, with  $n = 7$

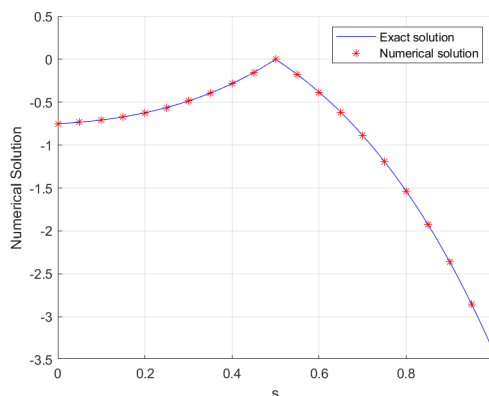


Fig. 6. Graphical comparison of  $\phi'$  vs  $\phi'_n$  for Example 2, with  $n = 7$

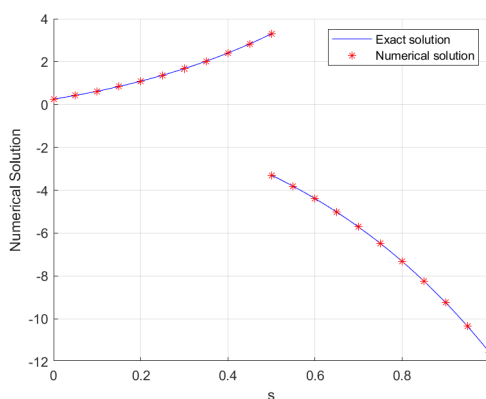


Fig. 7. Graphical comparison of  $\phi''$  vs  $\phi''_n$  for Example 2, with  $n = 7$

**Discussion:** We observe from the tables that the error function is almost null, especially for big values of the number  $n$ . Therefore, our proposed method is more effective whenever the degree of the polynomial (approximate solution)  $n$  is greater.

## 7. Conclusion

In this paper, the Galerkin method has been applied using the Chebyshev wavelets to approach the exact solution for a nonlinear Fredholm integro-differential equation of the second order with the initial conditions. The present method allowed us to reduce the equation into a nonlinear algebraic system, then we solve this system using the MATLAB tool. The numerical examples have proved the accuracy and validity of the proposed method.

## References

- [1] Thieme, H.R. (1977). A model for the spatial spread of an epidemic. *Journal of Mathematical Biology*, 4(2), 337-351.
- [2] Jerri, A.J. (1971). *Introduction to Integral Equations with Applications*. New York: Marcel Dekker.
- [3] Boltzmann, L. (1872). Further studies on the thermal equilibrium among gas-molecules. *Collected Works 1*, 275-370.
- [4] Balesku, R. (1963). *Statistical Mechanics of Charged Particle*. London: Wiley.
- [5] Fathy, M., El-Gamel, M., & El-Azab, M.S. (2014). Legendre-Galerkin method for the linear Fredholm integro-differential equations. *Applied Mathematics and Computation*, 243, 789-800.
- [6] Kumbinarasaiah, S., & Mundewadi, R.A. (2021). The new operational matrix of integration for the numerical solution of integro-differential equations via Hermite wavelet. *SeMA Journal*, 78(3), 367-384.
- [7] Erfanian, M., Gachpazan, M., & Beiglo, H. (2017). A new sequential approach for solving the integro-differential equation via Haar wavelet bases. *Computational Mathematics and Mathematical Physics*, 57(2), 297-305.
- [8] Dhawan, S., Machado, J.A.T., Brzeziński, D.W., & Osman, M.S. (2021). A Chebyshev wavelet collocation method for some types of differential problems. *Symmetry*, 13(4), 536.
- [9] Adibi, H., & Assari, P. (2010). Chebyshev wavelet method for numerical solution of Fredholm integral equations of the first kind. *Mathematical Problems in Engineering*, 5, ID 41449201.
- [10] Avudainayagam, A., & Vani, C. (2000). Wavelet-Galerkin method for integro-differential equations. *Applied Numerical Mathematics*, 32(3) 247-254.
- [11] Shiralashetti, S.C., & Kumbinarasaiah, S. (2019). Laguerre wavelets collocation method for the numerical solution of the Benjamina-Bona-Mohany equations. *Journal of Taibah University for Science*, 13(1), 9-15.
- [12] Bahuguna, D., Ujlayan, A., & Pandey, D.N. (2009). A comparative study of numerical methods for solving an integro-differential equation. *Computers & Mathematics with Applications*, 57(9), 1485-1493.
- [13] Biçer, G.G., Öztürk, Y., & Gülsu, M. (2018). Numerical approach for solving linear Fredholm integro-differential equation with piecewise intervals by Bernoulli polynomials. *International Journal of Computer Mathematics*, 85(10), 2100-2111.
- [14] Pittaluga, G., & Sacripante, L. (2009). An algorithm for solving Fredholm integro-differential equations. *Numerical Algorithms*, 50(2), 115-126.
- [15] Mahmoodi, Z., Rashidinia, J., & Babolian, E. (2013). B-Spline collocation method for linear and nonlinear Fredholm and Volterra integro-differential equations. *Applicable Analysis*, 92(9), 1787-1802.
- [16] Al-Khaled, K. (2002). On the rate of convergence for the Chebyshev series. *Missouri Journal of Mathematical Sciences*, 14(1), 4-10.